states. This can be seen, for example, as a finite zero-bias conductance in tunneling spectroscopy, and may be responsible for the excess quasi-particles identified in bulk measurements.

The above data give direct evidence of a strong source of low-energy quasi-particle excitations in BSCCO even at low temperatures. The observed QPSR characteristics, including breaking of particle-hole symmetry, diameter of  $2\xi_0$ , and LDOS decay as  $1/r^2$ , are all consistent with strong quasi-particle scattering from atomic-scale scatterers in a *d*-wave (but not *s*-wave) superconductor. Finally, this experiment demonstrates the possibility of using a single atomic-scale scatterer as a probe of high- $T_{\rm c}$  superconductivity. This technique might eventually be used, with specific elemental impurities deliberately doped into different atomic planes, to help unlock the secrets of the mechanism of high- $T_{c}$  superconductivity.

*Note added in proof*: Since the submission of this manuscript, the authors have become aware of a recent study on atomic-scale impurities (*32*).

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- 33. We acknowledge and thank A. Balatsky, M. Flatté, M. Franz, A. de Lozanne, K. Moler, J. Orenstein, R. E. Packard, and D. J. Scalapino for helpful conversations. Funded by NSF under grants DMR-9458015 and DMR-9623842, by the D. & L. Packard Foundation, and by the U.S. Department of Energy.

19 March 1999; accepted 19 May 1999

# High H<sub>2</sub> Uptake by Alkali-Doped Carbon Nanotubes Under Ambient Pressure and Moderate Temperatures

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Lithium- or potassium-doped carbon nanotubes can absorb  $\sim 20$  or  $\sim 14$  weight percent of hydrogen at moderate (200° to 400°C) or room temperatures, respectively, under ambient pressure. These values are greater than those of metal hydride and cryoadsorption systems. The hydrogen stored in the lithium-or potassium-doped carbon nanotubes can be released at higher temperatures, and the sorption-desorption cycle can be repeated with little decrease in the sorption capacity. The high hydrogen-uptake capacity of these systems may be derived from the special open-edged, layered structure of the carbon nanotubes made from methane, as well as the catalytic effect of alkali metals.

Hydrogen has been recognized as an ideal energy carrier, but to make it truly useful, enduser H<sub>2</sub> storage must be improved. In particular, high storage capacity is desirable when H<sub>2</sub> is used as the energy carrier in high-energy density rechargeable batteries and in H<sub>2</sub>/O<sub>2</sub> fuel cells. For these applications, metal hydridation is the existing method above room temperatures and below 20 to 40 atm of pressure, but these materials are heavy and expensive (1, 2). Cryoadsorption, in which activated carbon is often used as a sorbent, shows its advantages in the moderate size and weight of carbon, but suffers from the severe conditions (liquid nitrogen temperatures and 20 atm of pressure) required to hold the physically adsorbed  $H_{2}$  (3, 4). In any case, the  $H_2$  uptake by these systems is less than 6 weight  $\sqrt[6]{}$  (Table 1), far lower than that of gasoline or diesel (17.3 weight %). More recently, carbon nanotubes were reported to be a more effective material for H<sub>2</sub> uptake. Dollin et al. found that single wall carbon nanotube (SWNT) soots could absorb about 5 to 10 weight % of H<sub>2</sub> at 133 K and 300 torr (5). Chambers et al. observed that at 120 atm and room temperature, graphite nanofibers with herringbone structure could store 67 weight % of  $H_2(6)$ . Ye et al. used high-purity SWNT and obtained 8.25 weight % of H<sub>2</sub> adsorption at 80 K and 100 atm (7). All the above  $H_2$ -uptake systems require high pressure or subambient

temperatures, or both. Here we introduce a H<sub>2</sub> storage system that uses alkali metal–doped carbon nanotubes (CNTs) as sorbents and operates at ambient pressure and moderate temperatures. The H<sub>2</sub> uptake can achieve 20 weight % for Li-doped CNT at 653 K, or 14 weight % for K-doped CNT at room temperature. These values correspond to ~160 (for Li-doped CNT) or 112 kg of H<sub>2</sub>/m<sup>3</sup> (for K-doped CNT), respectively, and are comparable to those of gasoline and diesel.

The CNTs used in this study were made from catalytic decomposition of  $CH_4$  (8). After purification, almost all of the catalyst particles were removed. More than 90% of the product was in the form of multiwalled CNTs, and 70% was in the diameter range of 25 to 35 nm. The structure of a CNT is formed by the piling up of graphene sheets in the shape of circular cones with a hollow center. The doping of Li and K to the CNT was carried out by solid-state reactions between CNT and Li- or K-containing compounds, such as carbonates or nitrates. For comparison, Li- and K-doped graphite were prepared by the same procedures. The graphite sample was obtained from Merck with an average diameter of 50 µm. The specific surface area of CNT and graphite is 130 and 8.6 m<sup>2</sup>/g, respectively. The Li/C and K/C ratio of these alkali-doped carbon materials was about 1/15 as measured by x-ray photoelectron spectroscopy. The density of Lidoped carbon materials was  $\sim 0.9$  g/cm<sup>3</sup> for CNT and  $\sim 2.0$  g/cm<sup>3</sup> for graphite. Hydrogen

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**Fig. 1.** TGA profile of a reversible  $H_2$  sorptiondesorption cycle on (**A**) Li-doped CNT and (**B**) K-doped CNT. Samples were initially treated at 873 K for 1 hour. For Li-doped CNT, the sample was cooled to 300 K, and then heated again to 873 K. For K-doped CNT, the sample was cooled to room temperature (298 K), maintained at 298 K for 2 hours, and then heated to 773 K.

uptake was measured by thermogravimetery analysis (TGA), with purified  $H_2$  (>99.99%) as the purging gas. Hydrogen absorptiondesorption was confirmed by temperatureprogrammed desorption (TPD), with  $H_2$  being the only desorption product of the  $H_2$ saturated carbons. In situ Fourier transform infrared spectroscopy (FTIR) was applied to analyze the detailed mechanism of the process. All the above investigations were performed mainly on Li-doped samples because they are stable under ambient conditions.

Samples for TGA were initially heated in situ at 873 K for 1 hour in a flow of purified  $H_2$  (99.99%) to remove absorbed water and contaminants. The Li-doped samples were cooled from 873 to 300 K and then heated again to 873 K linearly (5° per minute). The  $H_2$  uptake began at temperatures around 773 K and ended at 423 K (Fig. 1A). When the temperature was increased again, the sample further increased its weight, reaching a maximum value corresponding to a total of 14.5 weight % H<sub>2</sub> at 673 K. Further increasing the temperature resulted in the H<sub>2</sub> desorption, and the sample weight returned to its original value at 773 K. In the above hydrogenationdehydrogenation cycle, the system did not reach the equilibrium (saturation) state.



**Fig. 2.** FTIR spectra of (a) CNT at 650 K, and (b to d) Li-doped CNT in the  $H_2$  stream at 873 K for 1 hour, 650 K for 1 hour, and 650 K more than 2 hours, respectively.

When the Li-doped CNT sample was maintained at 653 K for 2 hours, the  $H_2$  uptake reached 20 weight %  $H_2$  at ambient pressure.

The K-doped samples were cooled from 873 K to room temperature in a H<sub>2</sub> stream and maintained at room temperature for 2 hours. They were then heated again at 5° per minute to 773 K. A gradual weight increase was observed from 773 to 343 K (Fig. 1B). As the temperature was further decreased, a rapid weight increase occurred that reached  $\sim 14$  weight % (saturated value) after the system was maintained at room temperature for 2 hours. Desorption of H<sub>2</sub> took place as the temperature was raised. The desorption was rapid between 300 and 423 K and then became relatively gradual. It appears that there are at least two kinds of absorbed H<sub>2</sub> in K-doped CNT that desorb in temperature regimes.

Similar  $H_2$  uptake was observed for Li- and K-doped graphite, but the  $H_2$  uptake by alkalidoped graphite was only 35 to 70% of that by alkali-doped CNT (Table 1). It seems that  $H_2$ uptake capacity correlates with the structure of carbon sorbents. Compared with graphite, the carbon nanotubes used in this work consist of small-sized graphene sheets (50 nm or less) in the shape of a hollow circular cone and have much more open edge and greater interplanar distance (0.347 nm for CNT versus 0.335 nm for graphite). All these features favor the high  $H_2$  uptake by alkali-doped CNT.

Compared with 0.4 weight % or no H<sub>2</sub> absorption measured for the same carbon nanotubes or graphite sample without alkali doping (9), the high H<sub>2</sub>-uptake capacity may derive from the properties of alkali metals. Our in situ FTIR investigation has unambiguously shown that the H<sub>2</sub> uptake under the conditions used in this study (ambient pressure, moderate temperatures, presence of alkali metals) is actually a dissociative hydrogenation of the carbon sorbents. When H<sub>2</sub> uptake was not evident at temperatures above 773 K, there was only weak Li-H vibration at  $\sim$ 1420 cm<sup>-1</sup> (Fig. 2, curve b). As the temperature was reduced to 653 K, the Li-H vibration became more intense and a weak band, which is due to C-H stretching, developed between 3000 and 3350  $cm^{-1}$ (curve c). When the sample was kept in contact with H<sub>2</sub> at 653 K for hours, the C-H band appeared gradually and broadened, whereas no significant change in the Li-H vibration was observed (Fig. 2, curve d). The infrared absorption corresponding to C-H stretching in alkanes occurs in the range of 2840 to 3000  $\text{cm}^{-1}$ , whereas in cyclic alkanes or alkenes absorption it is between 3000 and 3300 cm<sup>-1</sup> (10). This newly developed band at 2600 to 3400  $\text{cm}^{-1}$ therefore revealed the formation of CH, species, indicating that the adsorption resulted from the dissociative hydrogenation of carbon nanotubes or graphite. Upon heating to high temperatures at which the H<sub>2</sub> desorption took place, the C-H band correspondingly vanished.

The formation of the LiH species was observed by FTIR to occur before C-H formation, and it maintained constant intensity, whereas the C-H band increased with increasing  $H_2$  exposure. This suggests that Li may act as a catalytic active center for the  $H_2$ dissociative adsorption. The dissociated H atoms may spill over from Li sites to the carbon network of grapheen sheets, and finally become bonded to carbon atoms. Our x-ray

**Table 1.** Comparison of  $H_2$  storage properties of various systems.  $W_1$ :  $H_2$ -uptake capacity (weight %  $H_2$ );  $V_1$ :  $H_2$  density (g/liter).

System	T <sub>absorb</sub> (K)	P <sub>absorb</sub> (atm)	H <sub>2</sub> density		Energy density	
			W	V,	kWh/Kg	kWh/liter
CNT	298-773	1	0.4	3.2	0.133	0.106
_i-doped						
CNT	473~673	1	20.0	180	6.66	6.0
Graphite	473~673	1	14.0	280	4.66	9.32
<-doped						
ĊŃŢ	<313	1	14.0	12.6	4.66	4.2
Graphite	<313	1	5.0	60	1.66	2.0
FeTi-H	>263	25	<2	96	0.58	3.18
NiMg-H	>523	25	<4	81	1.05	2.69
Cryoadsorption	~77	20	~5	~20	1.66	0.67
sooctane/gasoline	>233	1	17.3	117	12.7	8.76

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diffraction measurement revealed the formation of Li<sub>2</sub>C<sub>2</sub> in the Li-doped CNT or graphite. We have also performed ultraviolet photoelectron spectroscopy (UPS) studies of valence band structure for both samples with and without Li. It was shown that the Li doping resulted in an extra half-filled electron-density-of-state containing the Fermi edge (11). The H<sub>2</sub>-dissociative adsorption on carbon is a slow activated process (12), with an activation energy corresponding to an above-zero-energy crossing between the di-H atoms and H<sub>2</sub> molecular potential curves. Theoretical band-structure calculation (13) has indicated that the half-filled Fermi level band created by the Li doping can overlap strongly with the unoccupied antibonding H<sub>2</sub>  $(1s^2)^*$  orbital, which to a large extent reduces the energy barrier for H<sub>2</sub> dissociation. We can therefore observe the high H<sub>2</sub> uptake resulting from Li doping.

The H2-rechargeability of Li- and Kdoped samples was tested by TGA. For Lidoped samples (CNT and graphite), the saturated H<sub>2</sub> uptake was measured at 653 K after each complete desorption at 823 K, whereas for K-doped carbon materials it was measured at 298 K after each run of desorption at 773 K. The results show that after more than 20 cycles of absorption-desorption, the capacities of H<sub>2</sub> uptake are reduced by less than 10% for both systems. High H<sub>2</sub> pressure was shown to favor the H<sub>2</sub> absorption, which is expected because H<sub>2</sub> uptake is a volumereducing process. TPD measurements have demonstrated that the Li-doped CNT exposed to  $H_2$  at 10 atm for 1 hour can store the same amount of H<sub>2</sub> as those systems at ambient pressure for 2 hours.

Although K-doped carbon samples can absorb  $H_2$  at lower temperature than Lidoped samples, Li-doped carbon materials are chemically more stable than K-doped carbon materials. They can maintain  $H_2$  uptake capability even after being heated in air at 373 K for hours, and no flame resulted even when the samples were exposed to air at 673 K after  $H_2$  had been absorbed. On the other hand, K-doped CNT can be oxidized rapidly and even cause fire shortly after being exposed to air at room temperature. Nevertheless, both systems may find wide applications in the near future.

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12 January 1999; accepted 1 June 1999

## Regulation of NMDA Receptors by an Associated Phosphatase-Kinase Signaling Complex

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Regulation of *N*-methyl-D-aspartate (NMDA) receptor activity by kinases and phosphatases contributes to the modulation of synaptic transmission. Targeting of these enzymes near the substrate is proposed to enhance phosphorylation-dependent modulation. Yotiao, an NMDA receptor-associated protein, bound the type I protein phosphatase (PP1) and the adenosine 3',5'-monophosphate (cAMP)-dependent protein kinase (PKA) holoenzyme. Anchored PP1 was active, limiting channel activity, whereas PKA activation overcame constitutive PP1 activity and conferred rapid enhancement of NMDA receptor currents. Hence, yotiao is a scaffold protein that physically attaches PP1 and PKA to NMDA receptors to regulate channel activity.

The molecular organization of the postsynaptic density (PSD) is thought to be essential for the fidelity and precision of synaptic signaling events. Clustering and immobilization of neurotransmitter receptors and ion channels is maintained by an intricate system of proteinprotein interactions (1). For example, NMDA receptors are clustered and coupled to the cytoskeleton through association with PDZ domain-containing proteins,  $\alpha$ -actinin, and neurofilaments (2). Many signaling pathways converge on the NMDA receptor (3), allowing the regulation of channel activity in response to the generation of second messengers such as Ca<sup>2+</sup> and cAMP (4, 5). PKA and PP1 activities modulate NMDA receptor function and appear to act in opposition to each other (5, 6). Individual targeting or anchoring proteins such as AKAP79 and spinophillin localize the kinase and phosphatase at the PSD (7, 8).

A two-hybrid screen for proteins that bind

the NMDA receptor subunit isoform NR1A identified a protein called votiao that interacts with the COOH-terminal C1 exon cassette of the ion channel (9). We isolated cDNAs encoding fragments of votiao by an interaction cloning strategy to identify A-kinase anchoring proteins (AKAPs) (10) and confirmed that the protein bound NR1A (11). Expression of fulllength votiao fused to green fluorescent protein (GFP) in HEK 293 cells (12) resulted in detection of a  $\sim$ 210-kD protein that bound the type II regulatory subunit of PKA (RII), as assessed by overlay assay (Fig. 1A). Immunoprecipitations with antiserum to yotiao from brain extracts also isolated an RII binding protein and were enriched by a factor of  $10.5 \pm 2$  (n = 3) for PKA catalytic subunit activity (Fig. 1, B and C). We mapped the RII binding site to a region between residues 1229 and 1480 by screening a family of recombinant yotiao fragments (Fig. 1D) in the overlay assay (Fig. 1, E and F). Residues 1440 to 1457 appear to include the principal determinants for RII interaction because a synthetic peptide encompassing this region blocked RII binding (Fig. 1F). These findings indicate that votiao functions to anchor PKA to NMDA receptors.

Because PP1 activity participates in the regulation of NMDA receptors ( $\delta$ ), we conducted experiments to address whether the

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