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## A Post-Stishovite SiO<sub>2</sub> Polymorph in the Meteorite Shergotty: Implications for Impact Events

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Transmission electron microscopy and electron diffraction show that the martian meteorite Shergotty, a shocked achondrite, contains a dense orthorhombic SiO<sub>2</sub> phase similar to post-stishovite SiO<sub>2</sub> with the  $\alpha$ -PbO<sub>2</sub> structure. If an SiO<sub>2</sub> mineral exists in Earth's lower mantle, it would probably occur in a poststishovite SiO<sub>2</sub> structure. The presence of such a high-density polymorph in a shocked sample indicates that post-stishovite SiO<sub>2</sub> structures may be used as indicators of extreme shock pressures.

The SiO<sub>2</sub> polymorph stishovite forms at high pressure (1), and has been found in highly shocked rocks (2). The existence and stability of post-stishovite silica (SiO<sub>2</sub> polymorphs denser than stishovite) have been the subject of many recent studies (3-14). Post-stishovite silica plus magnesiowüstite might be stable relative to (Mg.Fe)SiO<sub>3</sub>-perovskite in Earth's lower mantle (13, 15–17). Such an SiO<sub>2</sub> phase in heavily shocked samples would also provide a tool for understanding impact processes. Here we describe a post-stishovite SiO<sub>2</sub> phase in the martian meteorite Shergotty.

Shergotty is heavily shocked; it contains maskelynite (18–20), which was long thought to be diaplectic plagioclase glass that formed by a solid-state transformation (21), but appears to have quenched from a shock-induced melt (22). Shergotty also contains large silica grains (>150 mm) that were previously interpreted to be birefringent shocked quartz with planar deformation features (21). Most are wedge-shaped, which is typical of  $\beta$ -tridymite but atypical of quartz. Like many of the maskelynite grains (22), the silica grains lack shock-induced

intragranular fractures. Many of the SiO<sub>2</sub> grains are surrounded by radiating cracks (Fig. 1A). These cracks are similar to those formed around coesite in high-pressure metamorphic rocks (23)and indicate that the volume of the silica phase increased greatly with decompression. Expansion must have occurred when the maskelynite was solid, because the cracks also cut through it (Fig. 1B). The lamellar texture appears as two sets of lamellae of different brightness in fieldemission scanning electron microscopy images recorded in back-scattered-electron (BSE) mode (Fig. 1B). Electron microprobe (EMP) analyses of the largest lamellae (up to several micrometers wide), using a slightly defocused electron beam, revealed that both the bright and dark lamellae are SiO<sub>2</sub> compounds that are identical in composition and have minor concentrations of Na<sub>2</sub>O (0.4% by weight) and Al<sub>2</sub>O<sub>3</sub> (1.12%). BSE imaging with a field-emission scanning electron microscope (FESEM) at low accelerating voltages revealed that the original silica grains consist of a mosaic of many individual domains (10 to 50 µm in diameter), each with a distinct pattern of intersecting thin (<300 nm) lamellae (Fig. 1B).

We investigated the crystallinity and structure of  $SiO_2$  phases using laser Raman microprobe spectroscopy, transmission electron microscopy (TEM), and selected-area electron diffraction (SAED). Three zone-axis SAED patterns (Fig. 2) were collected from a domain of crystalline material. These patterns all show sharp intensity maxima and diffuse rings, corresponding to diffraction from the crystalline material and diffuse scattering from amorphous material, respectively. Starting from zone axis 1 (Fig. 2A), the crystal was tilted 21° about the vertical diffraction vector (with d spacing = 4.54 Å) to reach zone axis 2 (Fig. 2B). The reflection corresponding to 4.54 Å disappeared as the sample was tilted out of zone axis 1 and remained absent in zone axis 2, indicating that it is a forbidden reflection that occurs by dynamical diffraction in zone axis 1. Zone axis 3 (Fig. 2C) was reached by tilting the crystal 37° from zone axis 1 about the horizontal diffraction vector (d = 3.22 Å). In this case, the d = 3.22 Åreflection remained excited and therefore is not forbidden. The three zone axis patterns contain seven diffraction vectors that can be used to constrain the structure of this SiO<sub>2</sub> polymorph.

The diffraction data (Fig. 2) cannot be indexed using the structures of SiO<sub>2</sub> polymorphs that are stable at low pressures. Although tridymite is a likely SiO<sub>2</sub> phase in Shergotty before shock metamorphism, the diffraction data are inconsistent with all tridymite polymorphs (24). Similarly, the diffraction data are inconsistent with coesite and stishovite. The data are also inconsistent with known post-stishovite structures, including the CaCl<sub>2</sub>, baddelyite, and  $\alpha$ -PbO<sub>2</sub>, and modified baddelyite (Pnc2) structures of SiO<sub>2</sub>. The diffraction data from zone axis 1 fit a post-stishovite structure with space group Pbcn that is similar to the  $\alpha$ -PbO<sub>2</sub> structure (25-27), but the d spacings from zone axes 2 and 3 fit less well to this  $\alpha$ -PbO<sub>2</sub>-like structure (27). Our data also nearly fit a  $Pca2_1$  structure calculated to be a quench phase from post-stishovite SiO<sub>2</sub> at 50 GPa and a 10-GPa differential stress (28). The cell parameters of Pca2, and Pbcn structures are nearly the same, but they do not match the 1.97 Å and 3.41 Å d spacings of our data (Table 1).

Because the symmetries of our diffraction patterns are consistent with an orthorhombic lattice and the reflections can be indexed to *Pbcn* and *Pca2*<sub>1</sub> structures (27, 28), we refined new orthorhombic cell parameters to fit our data (Table 1). All our *d*-space data fit the refined orthorhombic unit cells within 1% (Ta-

**Table 1.** *d*-space data obtained from the three zone-axis diffraction patterns as compared to calculated *d* spacings for *Pbcn* and *Pca2*<sub>1</sub> SiO<sub>2</sub> structures (*27, 28*) and the refined unit cells based on the *Pbcn* structure. Miller indices (*hkl*) are also given for each structure, with an asterisk indicating the reflections that are forbidden

SAED	Pbcn		Pca2 <sub>1</sub>		Refined Pbcn	
4.54*	001*	4.50	100*	4.49	001*	4.55
3.41	011*	3.25	110	3.25	011*	3.40
3.22	110	3.17	011*	3.17	110	3.23
3.09	101*	3.11	101*	3.11	101*	3.07
2.62	111	2.59	111	2.59	111	2.63
2.28	002	2.25	200	2.25	002	2.28
1.97	121	1.88	121	1.88	121	1.97

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Fig. 1. BSE-mode SEM images of SiO<sub>2</sub> and maskelynite. (A) An SiO<sub>2</sub> grain and maskelynite are surrounded by a fractured clinopyroxene (cpx) where the fractures radiate more than 500  $\mu m$  from the SiO<sub>2</sub>. (B) FESEM image of a triangular SiO, grain showing the tweedlike internal microstructure with some nearly orthogonal lamellae and fractures radiating outward into the surrounding maskelynite.



ble 1). Our diffraction data also fit both refined lattices in terms of interplanar angles and angles between zone axes. Many forbidden reflections that are expected for the *Pbcn* and *Pca2*<sub>1</sub> space groups (Table 1) are present in our diffraction patterns as a result of dynamical diffraction effects. However, the 3.22 Å reflection (011 in *Pca2*<sub>1</sub>) is not extinct as it would be for *Pca2*<sub>1</sub>, *a* = (0kl, l = 2n). A monoclinic structure (*P2*<sub>1</sub>, *a* =



**Fig. 2.** SAED patterns of three zone axes from a small cluster of crystalline  $SiO_2$ . The diffraction vectors are marked by their corresponding *d* spacings. Although zones 1 and 2 share the same vertical direction in reciprocal space, the 4.54 Å reflection does not appear in zone 2.

4.58 Å, b = 4.17 Å, c = 5.13 Å, and  $\beta = 88.85^{\circ}$ ) has been calculated (28) that also fits our refined lattice parameters, but the extinct 4.54 Å reflection (100 in P2<sub>1</sub>) is not forbidden for this structure. Although our diffraction data are insufficient to determine the space group, they fit an orthorhombic unit cell ( $a = 4.16 \pm 0.03$  Å,  $b = 5.11 \pm 0.04$  Å,  $c = 4.55 \pm 0.01$  Å, cell volume  $V = 96.91 \pm 0.63$  Å<sup>3</sup>) and are consistent with the *Pbcn* space group of  $\alpha$ -PbO<sub>2</sub>. Assuming four formula units per cell as in *Pbcn*, the density is 4.12 gm/cm<sup>3</sup>.

The electron diffraction (Figs. 2 and 3) and bright-field imaging (Fig. 3) data also show that some of the SiO<sub>2</sub> is amorphous. This material produces a diffuse diffraction ring (Fig. 2), whereas the crystalline material produces discrete sharp reflections. In the brightfield image (Fig. 3), the strongly diffracting crystalline material appears darker than the amorphous SiO<sub>2</sub>. The crystalline grains are 50- to 200-nm-wide islands embedded in 50to 100-nm-wide amorphous veins (Fig. 3) that are similar in size to the fine features observed in the FESEM images (Fig. 1B). Many of the amorphous veins are parallel to (001) (using *Pbcn*) or subparallel to (110) (Fig. 3).

The tweedlike microstructure in the silica grains of Shergotty is not the result of shock-induced vitrification of quartz, as previously interpreted (21), but rather represents the post-shock vitrification of unstable SiO<sub>2</sub> polymorphs formed during shock. The orthogonal lamellae seen in BSE images correspond to amorphous veins and intergrowths of crystalline SiO<sub>2</sub> polymorphs (26). BSE images (Fig. 1B) also show domains with nearly constant lamellae orientations that may represent single crystals of high-pressure SiO<sub>2</sub> in a larger polycrystalline SiO<sub>2</sub> grain.

Laser Raman microprobe (LRM) analysis of optically visible  $SiO_2$  lamellae does not show evidence for the presence of known low- or high-pressure crystalline  $SiO_2$  polymorphs. Instead, the spectrum (Fig. 4) appears to be that of a high-density glass (29–31) similar to silica glass quenched from 30 GPa (31), but could

presumably also be that of a mixture of dense vitreous and crystalline  $SiO_2$ . There are several possible reasons for the lack of spectral features indicative of crystalline  $SiO_2$ : (i) the crystalline  $SiO_2$  polymorph may not have a (strong) Raman spectrum, (ii) the volume concentration of the crystalline phase within the irradiation volume is too small, (iii) the selection rules for the vibrational spectrum of the crystalline  $SiO_2$  phase break down because of its small crystallite size, or (iv) the crystalline material is unstable under the laser beam and vitrifies during acquisition of the Raman spectrum. The last explanation is likely because of the unstable nature of the SiO<sub>2</sub> in Shergotty (32).

The SiO<sub>2</sub> was extremely sensitive to irradiation damage during TEM and rapidly became amorphous even with low electron doses (32). The SiO<sub>2</sub> was much less stable than synthetic stishovite previously analyzed (33, 34). The extreme instability of the crystalline SiO<sub>2</sub> is consistent with the new orthorhombic structure being a dense silica polymorph. The observation that this highly unstable phase has survived the post-shock thermal history of Shergotty suggests that the post-shock temperatures of the meteorite were relatively low [see also (35, 36)].

Theoretical calculations and diamond-anvil experiments demonstrate that stishovite (P42/ mnm) transforms into the CaCl<sub>2</sub> structure by a displacive mechanism at pressures above about 45 GPa, but the SiO<sub>2</sub> polymorphs that are stable at higher pressures are controversial (10, 11). The following sequence of SiO<sub>2</sub> phase transitions has been proposed (10): stishovite ( $P4_2$ / mmm)  $\rightarrow$  CaCl<sub>2</sub> (*Pnnm*)  $\rightarrow$  modified baddelyite  $(Pnc2) \rightarrow pyrite (Pa3)$ . Dubrovinsky et al. (10) argue that the Pnc2 structure, intermediate between that of  $\alpha$ -PbO<sub>2</sub> and baddelyite, is the stable polymorph at pressure (P) > 80 GPA, whereas Teter *et al.* (11) argue that the  $\alpha$ -PbO<sub>2</sub> structure is stable above 80 GPa; the structures and total energies of Pnc2 and  $\alpha$ -PbO<sub>2</sub> SiO<sub>2</sub> phases are similar (11). Our SiO<sub>2</sub> structure is distinctly different from the  $\alpha$ -PbO<sub>2</sub> and Pnc2 structures but similar to a Pbcn structure produced

**Fig. 3.** Bright-field TEM image of SiO<sub>2</sub> showing dark islands of crystalline material (strongly diffracting) cut by numerous veins of lighter contrast amorphous material. The SAED pattern shows the [ $\overline{110}$ ] zone axis of the SiO<sub>2</sub> polymorph (based on the *Pbcn* setting) with a faint diffuse ring representing diffuse scattering by the amorphous veins. The amorphous veins occur predominantly parallel to (001) and subparallel to (110).



in shock experiments at 70 to 90 GPa (27) and a quench phase from diamond anvil experiments done without a pressure medium (6). The similarity between our orthorhombic phase, the phase created by Tsuchida and Yagi (6), and structures calculated to quench from 50 GPa and a differential stress of 10 GPa (28) suggest that the SiO<sub>2</sub> phase in Shergotty may have been stabilized by high differential stress rather than pressures in excess of 80 GPa.

The presence of Na<sub>2</sub>O and Al<sub>2</sub>O<sub>2</sub> in the SiO<sub>2</sub> grains, and their morphology, are consistent with the low-pressure precursor phase having been tridymite. The maskelynite and silica apparently melted during shock and solidified under high pressure. However, the abrupt contacts between angular SiO<sub>2</sub> grains and surrounding maskelynite (Fig. 1B), and the morphology of the silica grains (Fig. 1B), are inconsistent with both phases being liquids simultaneously. We have observed limited mixing of SiO<sub>2</sub> and maskelynite along some grain margins as well as textures suggesting injection of SiO<sub>2</sub> into fractures. We conclude that the poststishovite SiO<sub>2</sub> formed primarily by a solidstate transformation and to a lesser extent by melting and subsequent crystallization.

The presence of post-stishovite  $SiO_2$  in Shergotty implies that the inferred 29-GPa



Fig. 4. Bulk Raman analysis of 1- to  $2-\mu m$  areas in both light and dark lamellae showed no indication of crystalline material.

shock pressure (21) is too low, and the shock pressure probably exceeded the stability range of stishovite (>45 GPa). The new polymorph cannot have formed through decompression of the CaCl<sub>2</sub> structure because the latter is unquenchable and inverts to stishovite upon decompression (8-11). Instead, it is likely to be a quench product from a post-stishovite structure such as  $\alpha$ -PbO<sub>2</sub>. The maximum shock pressure is uncertain. In some experiments extensive melting and extreme shock deformation take place in excess of 80 GPa (37), whereas experiments on chondritic samples up to 83 GPa show little melting (38-40). The disparity between the 29-GPa and the >80-GPa pressure implied by  $\alpha$ -PbO<sub>2</sub>-like SiO<sub>2</sub> cannot be resolved until more is known about the stability of these phases and the role of differential stress in post-stishovite-SiO2 phase transitions.

The stability of (Mg,Fe)SiO<sub>3</sub>-perovskite in the deep lower mantle is dependent on the structures and free energies of the SiO<sub>2</sub> phases. If (Mg,Fe)SiO<sub>3</sub>-perovskite decomposes to SiO<sub>2</sub> plus magnesiowüstite in the lower mantle (16-17), the SiO<sub>2</sub> would have a post-stishovite structure. The present results confirm the existence of such a structure in a natural sample.

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