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- Substantially temperature-independent (±20%) susceptibilities lay between -2.6 and -1.2 (×10⁻³ electromagnetic units per mole) for the

Crystal Structure of the Polymer Electrolyte Poly(ethylene oxide)₃:LiCF₃SO₃

P. Lightfoot, M. A. Mehta, P. G. Bruce*

lonically conducting polymers (polymer electrolytes) are under intensive investigation because they form the basis of all solid-state lithium batteries, fuel cells, and electrochromic display devices, as well as being highly novel electrolytes. Little is known about the structures of the many crystalline complexes that form between poly(ethylene oxide) and a wide range of salts. The crystal structure is reported of the archetypal polymer electrolyte poly(ethylene oxide)₃:LiCF₃SO₃, which has been determined from powder x-ray diffraction data. The poly(ethylene oxide) (PEO) chain adopts a helical conformation parallel to the crystallographic *b* axis. The Li⁺ cation is coordinated by five oxygen atoms—three ether oxygens and one from each of two adjacent CF₃SO₃⁻ groups. Each CF₃SO₃⁻ in turn bridges two Li⁺ ions to form chains running parallel to and intertwined with the PEO chain. There are no interchain links between PEO chains, and the electrolyte can be regarded as an infinite columnar coordination complex.

Polymer electrolytes consist of salts, such as LiClO₄, that are dissolved in coordinating polymers, such as poly(ethylene oxide) $[(CH_2CH_2O)_n]$. They can exhibit ionic conductivities approaching that of many liquid electrolytes. Discovered in 1973 by Wright and co-workers (1), their potential as completely novel electrolytes with a solid solvent was recognized in 1978 by Armand and co-workers (2), who have extensively developed and exploited them. The ability to prepare polymer electrolytes in the form of thin (3 to 30 μ m) flexible films of infinite length and with a high ionic conductivity has resulted in the development of a new generation of electrochemical devices with a solid electrolyte and solid electrodes (3). Such all-solid-state cells have a number of advantages; they can be fabricated in a wide variety of shapes, the leakage of corrosive liquids is eliminated, and complete devices can be made by automated film processing. Among the devices under development are electrochromic displays, "smart" windows, and rechargeable lithium batteries. The last

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of these can store large quantities of energy per unit weight and volume. Major investment in the United States, Europe, and Japan is targeted at the development of such batteries for electric vehicle and conthree examples. No superconductivity was evident at 100 Oe and ≥1.4 K.

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sumer electronic applications (4).

Progress in the field of polymer electrolytes has been severely restricted by the lack of structural data. Part of the difficulty is determining the crystal structures of polymer:salt complexes by single-crystal methods (5, 6). We have established the crystal structure of PEO3:NaClO4 (7) from highresolution powder x-ray data. The archetypal polymer electrolyte, LiCF₃SO₃ dissolved in poly(ethylene oxide), is the most widely investigated system, both for technological applications and fundamental studies. Here we report the elucidation of the crystal structure of the important PEO₃:LiCF₃SO₃ complex from powder x-ray diffraction data (8). The use of lowtemperature data was critical to the development of a stable structural model through Rietveld refinement. The increase in quality of the high-angle data due to the lowering of thermal motion in such a "soft" molecular system is a significant factor when studying such systems.

Initial indexing of the room-temperature data was carried out with the program

Table 1. Refined atomic parameters for PEO₃:LiCF₃SO₃ at 298 K, space group $P2_1/a$, a = 16.768(2) Å, b = 8.613(1) Å, c = 10.070(1) Å, and $\beta = 121.02(1)^\circ$. Hydrogen atoms were included in the refinement but are not reported.

Atom	x	у	Ζ	U _{iso} * (Ų)
O(1)	0.206(2)	0.110(3)	0.090(3)	0.026(7)
C(1)	0.262(2)	0.091(3)	0.253(3)	0.026(7)
C(2)	0.366(2)	0.099(2)	0.298(2)	0.026(7)
O(2)	0.386(1)	0.246(2)	0.265(2)	0.026(7)
C(3)	0.486(1)	0.261(3)	0.341(3)	0.026(7)
C(4)	0.513(1)	0.426(3)	0.314(2)	0.026(7)
O(3)	0.455(1)	0.469(2)	0.158(2)	0.026(7)
C(5)	0.450(2)	0.631(2)	0.131(3)	0.026(7)
C(6)	0.387(2)	0.651(3)	-0.051(4)	0.026(7)
S(1)	0.2697(7)	0.124(2)	0.809(1)	0.084(9)
C(7)	0.328(1)	0.077(2)	0.703(2)	0.084(9)
O(4)	0.277(2)	-0.001(2)	0.903(2)	0.11(1)
O(5)	0.303(2)	0.264(2)	0.893(3)	0.11(1)
O(6)	0.172(1)	0.150(4)	0.693(2)	0.11(1)
F(1)	0.321(2)	0.204(2)	0.624(3)	0.24(1)
F(2)	0.284(2)	-0.031(2)	0.601(3)	0.24(1)
F(3)	0.413(1)	0.041(3)	0.798(2)	0.24(1)
Li(1)	0.156(5)	0.903(11)	0.906(9)	0.05

*Values constrained according to atom type. Li not refined.

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Fig. 1. Final observed (plus marks), calculated (solid line), and difference profiles (below) for the Rietveld refinement of PEO₃:LiCF₃SO₃ at 298 K.

TREOR (9). The first 20 observable lines were indexed to give a monoclinic cell a =10.064 Å, b = 8.613 Å, c = 14.441 Å, and $\beta =$ 95.65°, space group P2₁/n. This cell can be transformed to the alternative setting in P2₁/a: a = 16.77 Å, b = 8.613 Å, c =10.064 Å, and $\beta =$ 121.0°, which bears a resemblance to Chatani's cell for the PEO₃:NaSCN complex (6) (a = 16.83 Å, b =7.19 Å, c = 10.64 Å, and $\beta =$ 125.5°). On this basis, the structure of PEO₃:NaSCN was used to provide a partial starting model for the Rietveld refinement of the present structure, by initially using our 100 K data (10). Preliminary refinement then involved least-squares minimization of this partial model against the powder diffraction pattern, but limited to the range of bond distances and angles specified by the constraints, with the program GSAS (11). At this stage the fit to the raw diffraction data was sufficiently good to allow successive difference Fourier techniques to be applied with Rietveld refinement in order to determine the remainder of the structure (specifically, the location of the CF₃SO₃⁻ moiety), with soft constraints again being added to stabilize the geometry of the CF₃SO₃⁻ group as the model was developed [S–C = 1.82(1) Å, S–O =



Fig. 2. View of the PEO_3 :LiCF $_3SO_3$ structure along the chain axis (*b*).

1.42(1) Å, C-F = 1.32(1) Å, and tetrahedral angles of 109°]. The soft constraints effectively act as additional observations, which were necessary to improve the datato-parameter ratio and thus stabilize the refinement of such a complex structure from relatively poor powder diffraction data. The position of the Li+ ion was determined geometrically at a point within the PEO chain with distances from the oxygens that are typical for Li⁺-O coordination and then was allowed to refine freely. This procedure significantly improved the profile fit and yielded a simple geometry around the Li⁺ position (see below). Attempts to refine the Li⁺ ion in other positions were unsuccessful. In the final refinement cycles, isotropic temperature factors U_{iso} were given to all atoms, and these were grouped according to atom type (see Table 1). The temperature factors for Li and H were not refined.

The large values of U_{iso} for the atoms in the CF₃SO₃⁻ group may suggest one of three possibilities: (i) a site occupancy less than unity; (ii) dynamic (thermal) disorder of the CF₃SO₃ group; and (iii) static (positional) disorder. Our model naturally requires a stoichiometry PEO₃:LiCF₃SO₃, in agreement

Atoms	Distance (Å)	Atoms	Angle (degree)
O1_C1	1.42(1)	C6–O1–C1	108(1)
C1_C2	1.56(1)	O1–C1–C2	108(1)
C2–O2	1.40(1)	C1_C2_O2	110(1)
O2–C3	1.45(1)	C2_O2_C3	109(1)
C3–C4	1.55(1)	O2–C3–C4	111(1)
C4–O3	1.41(1)	C3–C4–O3	110(1)
O3–C5	1.41(1)	C4O3C5	115(1)
C5–C6	1.58(1)	O3–C5–C6	105(1)
C6–O1*	1.44(1)	C5–C6–O1*	106(1)
S1–O4	1.39(1)	01–Li1–O2	174(5)
S1–O5	1.42(1)	O1–Li1–O3	85(3)
S1–O6	1.46(1)	O1–Li1–O4	77(3)
S1–C7	1.83(1)	01–Li1–O5	83(3)
C7–F1	1.31(1)	O2-Li1-O3	94(4)
C7–F2	1.30(1)	02–Li1–O4	99(3)
C7–F3	1.28(1)	02–Li1–O5	104(4)
Li1–01	2.38(9)	O3–Li1–O4	134(5)
Li1–02	2.01(8)	O3–Li1–O5	107(4)
Li1–O3	1.72(7)	O4-Li1-O5	112(4)
Li1–O4	2.21(8)		
Li1–05	2.14(8)		

Table 2. Selected bond lengths and angles.

*Symmetry operator: (i) 1/2 - x, 1/2 + y, -z.

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Fig. 3. View of the structure along *c*. $CF_3SO_3^-$ groups are shaded. For clarity, coordination around one Li⁺ is shown in dashed lines.



Fig. 4. Local environment around Li⁺, showing coordination by three ether oxygens and oxygens from two neighboring triflate groups.

with the phase diagram study of Vallée, Besner, and Prud'homme (12). Several earlier reports have suggested stoichiometries of PEO_4 :LiCF₃SO₃ or $PEO_{3.5}$:LiCF₃SO₃ (13). This would require partial occupancy of the Li/CF3SO3 sites in our model. In order to explore the stoichiometry, we fixed the temperature factors for the Li/CF₃SO₃ species at 0.08 Å². After refinement, the occupancy value was 0.96(1), which sets an upper limit on the PEO:LiCF₃SO₃ ratio of 3.13(3). We can reasonably conclude that there is no significant deviation from the ideal 3:1 stoichiometry. The latter two possible explanations for the high thermal parameters were explored by carrying out low-temperature refinements. The relatively smooth variation of $U_{\rm iso}$ versus temperature observed for F and O suggests that this phenomenon is associated merely with thermal and not static disorder. The relatively large displacements (0.10 Å) for F even at low temperature are not surprising and may be expected from the very weak coordination of this part of the complex.

Results of the Rietveld refinement for the room-temperature structure are given in Tables 1 and 2. The final model consisted of 113 variables refined against 2600 data points and 88 soft constraints. The data range covered was $8^{\circ} < 2\theta < 60^{\circ}$ (Fig. 1). Final agreement factors were $R_{wp} = 3.84\%$ and $\chi^2 = 1.83$. Views of the structure along b and c are given in Figs. 2 and 3, respectively. The helical arrangement of the PEO chain within the unit cell is similar to that for the previously determined 3:1 complexes, NaSCN (5), NaI (6), and NaClO₄ (7). However, the disposition of the inorganic portion of the structure relative to this is significantly different than in the previous cases, reflecting the different size and shape of the anion and the smaller size of the cation. Focusing on one PEO chain (Figs. 2 and 3), the Li⁺ ions are located within it, forming a zigzag pattern along the axis of the helix (that is, b). Lithium is coordinated by five oxygens, composed of three PEO oxygens and one oxygen from each of two different CF₃SO₃⁻ groups (Figs. 3 and 4). The coordination is best described as trigonal bipyramidal, with the three PEO oxygens being almost coplanar. Each CF₃SO₃⁻ group associated with the chain bridges two adjacent Li⁺ ions, one oxygen being coordinated to each lithium, leaving the third oxygen free. There is little interchain interaction, the -CF3 groups being located in the space between the chains. The crystal structure demonstrates strikingly that these compounds are columnar coordination complexes, with long chains of Li⁺ ions being coordinated by ether oxygens and triflate anions (Fig. 2), the solid polymer results from chain entanglement rather than any ionic crosslinking. It is also evident from the crystal structure that replacement of the $-CF_3$ group with a bulkier, preferably asymmetric group, would inhibit crystallization. This is presumably the origin of the wide range of amorphicity exhibited by PEO:Li[(CF₃SO₂)₂N] and PEO:Li[(CF₃SO₂)₃C].

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Fire History and Climate Change in Giant Sequoia Groves

Thomas W. Swetnam

Fire scars in giant sequoia [Sequoiadendron giganteum (Lindley) Buchholz] were used to reconstruct the spatial and temporal pattern of surface fires that burned episodically through five groves during the past 2000 years. Comparisons with independent dendroclimatic reconstructions indicate that regionally synchronous fire occurrence was inversely related to yearly fluctuations in precipitation and directly related to decadal-to-centennial variations in temperature. Frequent small fires occurred during a warm period from about A.D. 1000 to 1300, and less frequent but more widespread fires occurred during cooler periods from about A.D. 500 to 1000 and after A.D. 1300. Regionally synchronous fire histories demonstrate the importance of climate in maintaining nonequilibrium conditions.

Despite the complexity inherent in local fire regimes, regional fire activity often oscillates in phase with year-to-year climatic variability. For example, the area burned annually across the southern United States tends to decrease during El Niño years and increase during La Niña years (1). This coupling between wildland fire and climate raises the possibility that intrinsic and stochastic factors, while contributing to local ecosystem heterogeneity, are overridden by

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regional climatic events and trends. Rapid, regional changes in vegetation may result from extreme climate-linked disturbances, such as catastrophic crown fires, whereas relatively slow vegetation changes may follow gradual, climate-driven shifts in surface-fire regimes (2). A broad range of spatial and temporal observations is therefore necessary to distinguish local- from regional-scale patterns and to encompass both high- and low-frequency changes in processes (3).

A commonly observed but rarely quantified phenomenon of disturbance regimes is

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