is no further turbidity development and (ii) the observed order of magnitude change in turbidity of early spring is reproduced. It is found that the springtime-enhanced UV dose under the ice is independent of these limits.

Using the data developed above, we have calculated the temporal development of UV radiation at 305 nm transmitted through sea ice without a snow cover. Two cases are illustrated in Fig. 1. The lower curve shows the decrease in transmitted UV light that accompanies the increase in ice turbidity associated with the springtime warming. The upper curve was generated assuming the changing atmospheric ozone content calculated in (1). Thus, Fig. 1 demonstrates a 20-fold increase in under-ice UV radiance in early October resulting from the coincidence of the presence of the ozone hole and the period of relatively high transparency for sea ice. This clearly has implications for organisms living within and under the ice.

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Response: The response of organisms to enhanced levels of ultraviolet radiation depends on numerous factors, only one of which involves atmospheric radiative transfer. Trodahl and Buckley make the very important point that the transmission of Antarctic ice decreases as spring progresses. Since the "ozone hole" is primarily a phenomenon of early spring, this suggests that potential biological effects of the ozone depletion may be larger than otherwise anticipated. Trodahl and Buckley show that the "hole" of October 1987 was accompanied by an increase in radiation dose beneath the ice by a factor of 20 as compared with that in years before the appearance of the ozone depletion. Despite the percentage increase in irradiance beneath the ice, the absolute radiation level is still small, since the albedo of ice remains large. The change in ice transmission over time is cause for concern, although a central issue involves a comparison between the radiation doses and radia-

Ice Volcanism on Ariel

The report "Solid-state ice volcanism on the satellites of Uranus" by David G. Jankowski and Steven W. Squyres (1) proposes a novel emplacement mechanism for surface "lava" flows on Ariel and Miranda (on Ariel the "lava" is almost certainly a mixture of water and ammonia with perhaps additional components). Whereas terrestrial lava flows are a mixture of liquid and solids (crystals), Jankowski and Squyres propose that the flows on the Uranian satellites were entirely crystalline during emplacement. Existing models of lava flows are capable of accounting for the parabolic cross sections measured by Jankowski and Squyres. It is thus incumbent on the authors to demonstrate that their novel mechanism is really required.

For the past 20 years it has been clear to volcanologists working on terrestrial lava flows (2) that flowing lava behaves, not as a viscous fluid, but as a "Bingham" fluid with a well-defined yield stress. A Bingham fluid is one that responds elastically to applied shear stresses until the stresses exceed its yield strength, after which it flows as a viscous fluid. Bingham rheology characterizes a wide class of mixtures of liquids with solid particles, such as suspensions of clay in water, pigments in oil (paint), rock debris in mud, and crystals in melted rock (lava). Studies of lava flows on Mount Etna, Hawaii, and the moon (2) support the idea that erupted lava is a Bingham fluid with a yield stress ranging from about 10^3 to 10^5 Pa, depending on silica content. The most characteristic aspect of such flows is the approximately parabolic profile of their margins (with suitable corrections when slopes exceed the angle of repose), which can be directly related to the Bingham yield strength, $Y_{\rm B}$. A formula valid for the profiles of lava flows, ice sheets, and debris flow lobes (all of which can be treated approximately as Bingham substances) relates the thickness of the flow's center h_0 to its horizontal width w:

$$Y_{\rm B} = \frac{\rho g h_0^2}{w}$$

where ρ is the density of the flow and g is the

tion tolerances of organisms beneath the ice. This topic clearly merits additional research.

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planet's surface acceleration of gravity (0.27 m/s^2 for Ariel).

Jankowski and Squyres use photoclinometry to measure the profiles of five probable "lava" flows on Ariel. As they show, these profiles can be adequately fit by parabolas (except, of course, where the slope of the parabola becomes too steep-mass movement and regolith processes act to gently taper the flow's edges). They then propose a model that treats the extruded material as a Newtonian viscous fluid spreading from a central vent until cooling raises the viscosity past the point where flow is possible. They show that this model predicts parabolic flow profiles as long as the "lava" is moving. They then assume that the profile of the flow does not change as it cools and stiffens and derive a viscosity from the distance the flow has traveled within the cooling time (estimated from flow thickness and the thermal diffusivity of water ice).

A much more natural explanation of the morphology of Ariel's "lava" flows is that the extruded material is a mixture of liquid and crystals and that the parabolic profiles are an expression of the Bingham yield strength of the mixture at the time of solidification. All information about the rheology during extrusion and flow is lost during solidification and cannot be recovered without additional information. In Table 1 we have used the parabolic fits of Jankowski and Squyres and the equation above to deduce the yield strength of the flow material. These yield strengths vary from 6.7 \times 10^3 Pa to 3.7×10^4 Pa, right in the midrange of terrestrial lavas. We do not

Table 1. Viscosities from (1) and Bingham yield stresses inferred from profiles in (1, figure 5).

Profile	Viscosity (Pa·s)	Bingham yield stress (Pa)
A B C D top D bottom	$\begin{array}{c} 4.5 \times 10^{15} \\ 3.5 \times 10^{15} \\ 9.0 \times 10^{14} \\ 1.1 \times 10^{15} \\ 9.0 \times 10^{14} \end{array}$	$\begin{array}{c} 3.7 \times 10^4 \\ 3.1 \times 10^4 \\ 1.4 \times 10^4 \\ 1.4 \times 10^4 \\ 6.7 \times 10^3 \end{array}$