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Superheavy Elements: An Early Solar System Upper Limit for Elements 107 to 110

Abstract. The abundance of samarium-152 in the Santa Clara iron meteorite is found to be 108×10^7 atoms per gram. This quantity, if attributed to fission of a superheavy element with atomic number 107 to 109, limits the amount of superheavy elements in the early solar system to 1.7×10^{-5} times the abundance of uranium-238. For element 110, the limit is 3.4×10^{-5}

Superheavy elements (SHE's) are elements with atomic numbers greater than 106 which may have been produced in stars by the same processes that produced thorium, uranium, and plutonium. Myers and Swiatecki (1) suggested that a region of relative stability, which would permit some of these elements to have appreciable half-lives, existed near atomic number Z = 114 and neutron number N = 184. More recent calculations, in which alpha and beta decay were considered in addition to spontaneous fission, suggest that the most stable nuclide is $^{294}110$, with a half-life $(t_{1/2})$ of 10^5 to 10^9 years (2). Numerous attempts have been made to synthesize SHE's in the laboratory and to search for them in nature. The possibility of an SHE with $t_{1/2} \approx 10^9$ years has led many investigators to search for extant SHE's; claims have been made for their discovery in halos in mica (3), in meteorites (4), and in hot springs (5), but these claims have been criticized (6). Other investigations are discussed in several reviews (7); there is currently no strong evidence for extant SHE's. Anders and co-workers (8, 9) claimed evidence for an extinct SHE with Z = 115 (or 114 or 113) in the Allende meteorite, based on xenon isotopic abundances that appear to be due to fission of an unknown transuranic nuclide. This result is difficult to understand in terms of calculated half-lives for SCIENCE, VOL. 214, 16 OCTOBER 1981

these nuclides, all of which are less than 1 year, based on the more recent calculations (2). However, the calculated halflives are uncertain by many factors of 10 (10).

We designed an experiment to look specifically at element 110, the only one expected to have a half-life long enough to have survived until the formation of the solar system. We chose an iron meteorite, Santa Clara, for our study because elements 107 through 110 are expected to be siderophilic, that is, concentrated in metallic phases, whereas the rare earth fission decay products are not. Santa Clara was chosen because it contains evidence for decay of extinct ¹⁰⁷Pd $(t_{1/2} = 6.5 \times 10^6 \text{ years})$ (11). Because this meteorite apparently formed very early after nucleosynthesis, it may contain evidence for SHE's with half-lives as short as a few million years. In order to explain the applicability of meteorites in addressing this problem, we will describe the mechanism by which iron meteorites are thought to have formed.

Initially, solid grains of metals and silicates condensed out of a hot solar nebula. These grains accreted to form larger bodies that eventually melted. The metal grains sank and coalesced to form an iron-nickel core with a silicate mantle. Later collisions with other bodies fragmented the object. A sample of this core was eventually perturbed into an earth-

crossing orbit, and fell as a meteorite.

By considering the thermodynamic properties of the elements, one can calculate the fraction of an element condensed into the grains as a function of temperature. Santa Clara is a typical member of the group IVB irons, which, on the basis of several elements, give a self-consistent grain condensation temperature of 1270 K at a pressure of 10^{-5} atm (12). At this temperature Fe and Ni are about 10 percent condensed, while Re, Os, Ir, and Pt, congeners of SHE's 107 to 110, are totally condensed. This work is mainly aimed at SHE 110, the only element expected to have a long half-life, but because the calculated halflives for fission are so uncertain, elements 107 to 109 are also considered.

To consider the properties of the SHE's in detail and to calculate the fraction condensed, we require estimates of their enthalpy (ΔH_v) and entropy $(\Delta S_{\rm v})$ of vaporization. Following other investigators (8, 13), we extrapolate values from congeners in other rows of the periodic table. The values of $\Delta H_{\rm v}$ are plotted in Fig. 1 (14). The extrapolation for elements 107 to 109 appears straightforward, with a clear monotonic trend toward increasing $\Delta H_{\rm v}$ with increasing period. For element 110 the extrapolation is less certain. We arbitrarily chose a value as indicated, but any value from 100 to about 180 kcal is plausible. For comparison, the value chosen by Anders et al. (8) for element 111 in a similar situation is indicated. Following the procedures of Kelly and Larimer (12), we calculated the fraction of each SHE condensed at 1270 K and 10^{-5} atm. Even the most volatile element, element 110, is 99.98 percent condensed under these conditions. Deviations from ideal solid solution, if similar to Pt in Ni (15), would increase the amount of element 110 condensed. Because the extrapolation of $\Delta H_{\rm v}$ for element 110 is uncertain, we will be conservative and assume that it is only 50 percent condensed; this is the observed fraction condensed of Pd, the most volatile congener of element 110 (16).

Abundances of the rare earth elements were determined by neutron activation analysis with radiochemical separation (17). The results are given in Table 1. Neutron activation analysis can only determine the quantity of the target nuclide; except for monoisotopic elements, the technique generally does not directly determine the total amount of the element present. The assumption of normal isotopic abundance to infer the total quantity of an element may not be appropriate in this instance, since significant amounts of several of the isotopes may be made by spontaneous fission.

In order to estimate the amount of SHE that could account for the observed rare earth element abundances, we need to know the fission yield of the SHE in the rare earth mass region. Steinberg and Wilkins (18) developed a theoretical model for mass distribution in fission. Their model explains the observation that as the mass of the fissioning nuclide increases, the heavy mass peak remains fixed at 88 neutrons (around mass 142), while the mass of the lighter fragment increases (19). They calculated the expected mass distribution for fission of ²⁹⁸114 and showed that the stability of the 88-neutron mass peak is preserved, and the other fragment, now heavier, has a mass distribution with a broad peak near mass 157. Element ²⁹⁴110 is expected to have a heavy mass peak near mass 153. Allowing for several neutrons puts the peak in the distribution slightly lower than the masses examined in this work, but as the peaks are quite flat, with a full width at half-maximum of about 20 mass units, the fission yield for ¹⁵²Sm is expected to be near the maximum. We assume a fission yield of 4 percent for ¹⁵²Sm (20).

To calculate a limit, we assume that all of the ¹⁵²Sm is due to SHE decay. The observed ¹⁵²Sm abundance implies a limit to the SHE abundance of 2.7×10^{10} atoms per gram. If the SHE totally condensed into Santa Clara as discussed above, the SHE/Ir ratio in Santa Clara will be equal to the ratio in the early solar system (21). Taking the observed Ir concentration in Santa Clara (22) and the solar Ir/U ratio (23) yields a limit to the solar SHE/238U ratio (atomic) of 1.7×10^{-5} . If the SHE were so volatile that it was only partially condensed, the limit would be higher. The assumption of 50 percent condensation for element 110 discussed earlier sets an upper limit to the $110^{/238}$ U ratio of 3.4×10^{-5} . Other estimates could be made on the basis of the other nuclides in Table 1, but the higher abundances would not further constrain the limit of SHE in Santa Clara. The nuclide with the lowest abundance, ¹⁷⁶Lu, is shielded from fission.

Runcorn (24) suggested that siderophile SHE's with half-lives of $\sim 10^8$ years were present in the early moon. The energetic decay of these elements supposedly melted the lunar interior and powered a lunar magnetic dynamo. Runcorn argued that remanent magnetism observed in lunar rocks was thermoremanent magnetism produced by cooling the rocks in a strong (~ 1 gauss) dipole field. He estimated that 3×10^{13} atoms Table 1. Abundances of rare earth elements in the Santa Clara meteorite.

| Nu- clide | Atoms of nuclide per gram* (× 10 ⁻⁷) | Grams of element per $gram^{\dagger}$ $(\times 10^{12})$ | Frac- tion of cosmic abun- dance \ddagger (× 10 ⁶) |
|-------------------|---|---|---|
| ¹⁵² Sm | 108 ± 18 | 1.0 | 0.23 |
| ¹⁵¹ Eu | 420 ± 100 | 2.0 | 1.3 |
| ¹⁵⁹ Tb | <340 | < 0.9 | < 0.8 |
| ¹⁶⁵ Ho | <100 | < 0.3 | < 0.16 |
| ¹⁷⁴ Yb | <400 | <3.0 | < 0.7 |
| ¹⁷⁶ Lu | 3.9 ± 2.4 | 0.44 | 0.46 |

*Uncertainties are 1 standard deviation (S.D.) based on counting statistics. Upper limits are 3 S.D. †Assuming normal isotopic composition. Based on 18 µg of iridium per gram in Santa Clara (26) and cosmic ratios of rare earth elements to iridium (27).

per gram would be required to power the dynamo. Our upper limit is about a factor of 10^3 less than this, suggesting that a lunar dynamo powered by SHE's 107 to 110 cannot be the source of lunar remanent magnetism.

The results of this work can be used to set limits on the SHE production rate or half-life (but not both). The recent discoveries of evidence for in situ decay of ²⁶Al $(t_{1/2} = 7.3 \times 10^5 \text{ years})$ (25) and ¹⁰⁷Pd (11) indicate that meteorites formed relatively soon (a few million



Fig. 1. The enthalpy of vaporization of the congeners of elements 107 to 109 increases monotonically with row in the periodic table. The extrapolation for element 110 is less certain, but as long as $\Delta H_{\rm v}$ of 110 is greater than about 90 kcal/mole, at least 50 percent of it will have condensed into the metal grains that formed Santa Clara. The extrapolated value for element 111 chosen by Anders and co-workers (8, 9) is shown for comparison.

years) after the last nucleosynthesis event. The nuclides above, however, are made by processes other than the rprocess, which is suggested for the formation of SHE. Two r-process nuclides, ¹²⁹I ($t_{1/2} = 1.6 \times 10^7$ years) and ²⁴⁴Pu $(t_{1/2} = 8.3 \times 10^7 \text{ years})$, are found in meteorites in amounts that suggest a period of $\sim 10^8$ years from the last *r*-process nucleosynthesis to meteorite formation (26). If we assume, for illustrative purposes only, that the SHE half-life is identical to that of 244 Pu, then the SHE/ 238 U limit of 1.7×10^{-5} found in this work, combined with a Pu/238U ratio of 0.015 (27), yields a limit to the SHE/ ²³⁸Pu production ratio of 0.0011. If, on the other hand, we assume that the production rates of ²⁴⁴Pu and the SHE are identical, we get a half-life of 1.2×10^7 years, assuming continuous nucleosynthesis and 10^8 years of free decay.

It has been suggested that neutroninduced fission may terminate the rprocess before it reaches SHE masses (28), but Schramm and Fiset (29) noted that r-process synthesis of SHE's was strongly dependent on a surface asymmetry parameter. They suggested that reasonable values would permit synthesis of SHE's. Steinberg and Wilkins (18) noted that the peak in cosmic abundances around mass 165, most likely due to fission at the termination of r-process nucleosynthesis, requires that the rprocess terminate beyond mass 300. It is not clear, however, that these high mass species can attain the region of stability through beta decay, because spontaneous fission may effectively compete with beta decay.

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- 16. Based on observed Pd (11) and Ir (22) contents
- and the solar Pd/Ir ratio (23).
 17. The sample was etched with dilute aqua regia before irradiation at a flux of 1.5 × 10¹⁴ cm⁻² sec⁻¹ for 24 hours. After irradiation it was again table (0.4 monet of the more manual) and etched (0.4 percent of the mass removed) and then dissolved with rare earth carriers. The rare earth fraction was separated as a group, counted on a Ge(Li) detector, and compared with chemi-cally prepared standards. Radiochemical yield was determined by reactivation of mixed carrier solution. The nuclides ¹³⁹La, ¹⁴⁰Ce, and ¹⁴⁶Nd, normally measured by neutron activation, could not be determined because their (n,γ) activation products had a significant contribution from neutron-induced fission of ²³⁵U.
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- We thank R. Frazier and D. Rollins for assist-ance, A. Meyer and the staff of the Missouri Research Reactor Facility for performing irra-diations, R. Reedy for helpful discussions, and C. Lewis and C. Moore of Arizona State Uni-versity for the sample of Santa Clara. This work was supported in part by NASA grant NSG-7436. The gamma-Ray Analysis Facility, estab-lished by NSF grant EAR 78-02708. Present address: Department of Earth and Plan-etary Sciences, Massachusetts Institute of Tech-nology. Cambridge 02139.
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Spectroscopic Detection of Stratospheric Hydrogen Cyanide

Abstract. A number of features have been identified as absorption lines of hydrogen cyanide in infrared spectra of stratospheric absorption obtained from a high-altitude aircraft. Column amounts of stratospheric hydrogen cyanide have been derived from spectra recorded on eight flights. The average vertical column amount above 12 kilometers is $7.1 \pm 0.8 \times 10^{14}$ molecules per square centimeter, corresponding to an average mixing ratio of 170 parts per trillion by volume.

We have analyzed high-resolution solar spectra at wavelengths near 3 μ m taken with a spectrometer aboard a Sabreliner jet aircraft flying in the lower stratosphere. This spectral region contains the strong v_3 band of HCN, with band center at 3311.48 cm⁻¹. A number of weak lines at the position of the HCN lines were identified. This spectral region also contains lines of H₂O and N₂O as well as some unidentified lines. The HCN lines are weak, so that the signalto-noise ratio in a single spectrum does not permit identification of HCN, but by averaging a large number of individual scans and observing lines from P_{16} to R_{16} we have made a positive identification. To our knowledge, HCN has not been observed previously in the stratosphere or unpolluted troposphere.

The method of obtaining and analyzing stratospheric spectra has been reported (1). A solar tracking system directs sunlight into a high-resolution (0.06 cm^{-1} full width at half-maximum, apodized) Fourier transform spectrometer. The instrumentation is flown aboard the Sabreliner at times near sunset or sunrise to enhance the absorption path length.

Figure 1 shows a portion of the measured spectrum in the region containing the P branch of the ν_3 band of HCN. The HCN lines and major lines due to water are indicated. This spectral region has not been extensively studied in the past. To establish the presence of HCN absorption in the stratospheric spectrum, a wider portion of the spectrum was examined for HCN ν_3 lines with the line positions of Rank et al. (2). Although

Table 1. Lines in the ν_3 band of HCN used for analysis. The line positions are from Rank et al, (2) and the intensities are based on measurements by Jaffe (δ); the band intensity inferred from the line intensities given by Jaffe agrees within 3 percent of that from Hyde and Hornig (7).

| Line | Position (cm ⁻¹) | Intensity at 296 K (cm ⁻¹ /column amount*) |
|--|--|---|
| $ \begin{array}{r} P_{13} \\ P_{12} \\ P_{11} \\ P_{10} \\ P_{9} \\ P_{8} \\ \end{array} $ | 3271.45 3274.65 3277.83 3280.99 3284.13 3287.25 | $\begin{array}{c} 2.822 \times 10^{-19} \\ 3.026 \times 10^{-19} \\ 3.219 \times 10^{-19} \\ 3.446 \times 10^{-19} \\ 3.625 \times 10^{-19} \\ 3.450 \times 10^{-19} \end{array}$ |

*Molecules per square centimeter.

some lines are overlapped by other absorbers, all unobscured features for quantum number $J \leq 16$ appear in the measured spectra. Recent long-path laboratory measurements (3) provided absorption spectra of O_3 and N_2O in this region. These laboratory measurements and line parameters in the Air Force Geophysics Laboratory (AFGL) compilation (4) have been used to identify most of the other features in this region and to select HCN features that are free from absorption by other molecules (5).

Six lines of HCN in the region 3270 to 3290 cm^{-1} were selected as free from overlap by other absorptions and sufficiently intense to be useful in determining amounts of HCN. The six lines used in the analysis are listed in Table 1 with their line intensities measured by Jaffee (6). The intensities are estimated to be accurate to \pm 20 percent or better, and they agree well with band intensities given by other investigators (7, 8). To determine the amount of absorber, synthetic spectra are calculated (9) to match the observed spectra. In the absence of better information, we assumed a uniform mixing ratio of HCN above the aircraft (10). Since the lines are all weak, any other distribution would produce the same absorption for the same line-ofsight amount of HCN; any errors introduced by this assumption in the conversion of line-of-sight amounts to column amounts are negligible compared to the noise. The temperature profile of the U.S. Standard Atmosphere was used. The upper curve in Fig. 1 is a calculated spectrum based on line parameters for H_2O , CO_2 , and N_2O from the AFGL line compilation (4) and the HCN line parameters in Table 1.

Table 2 gives details of the flight spectra that were analyzed to determine amounts of HCN. Flights were made at latitudes from 5° to 50°N in winter and summer. The errors in the determination of HCN column amounts are largely due to noise in the spectra; systematic errors due to assumptions of distribution, temperature effects, errors in line intensities or broadening coefficients, and so on are generally much smaller. To improve the signal-to-noise ratio, we averaged as many spectra as possible, even when this covered a significant range of air mass,